

7.5

Homogeneous Linear Systems with Constant coefficients.

$$\bar{x}' = A\bar{x}$$

$\det A \neq 0$, then $\bar{0}$ is the only critical point (CP)

If $\lim_{t \rightarrow \infty} \bar{x}(t) = \bar{0}$ then the

CP is asymptotically stable.

Stable if the solution stay close to $\bar{0}$ for all times.

Otherwise the CP is unstable

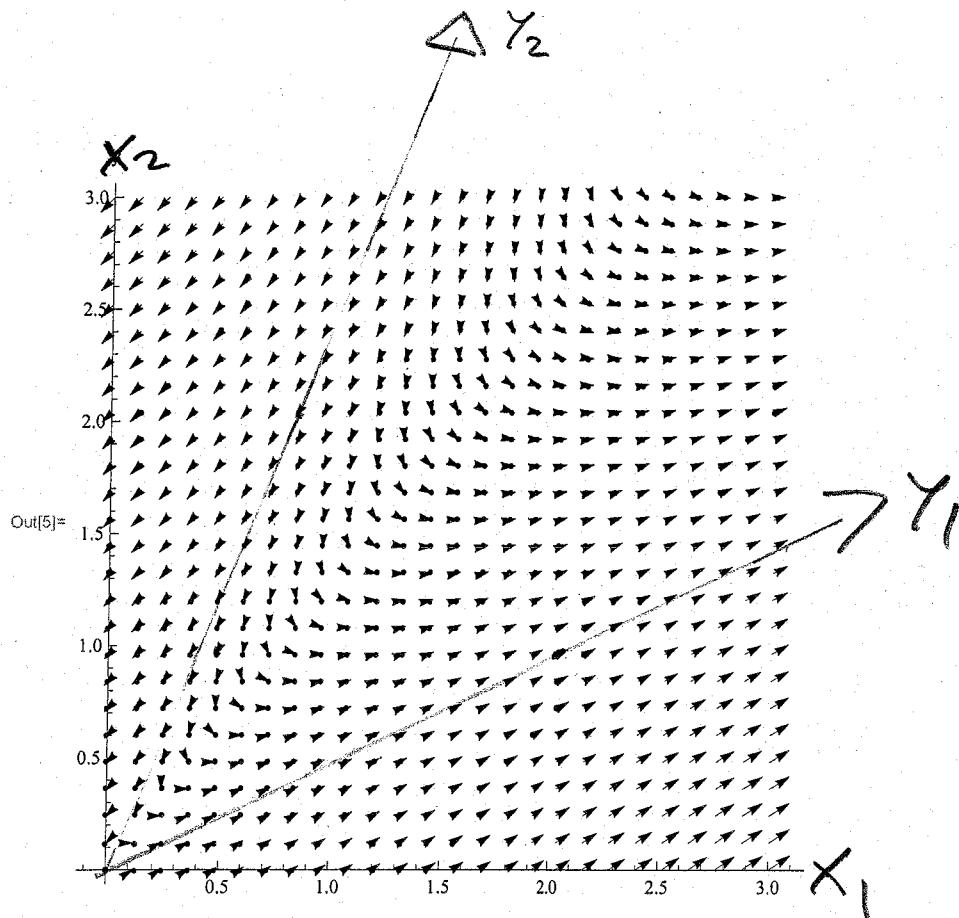
Ex) 7.5.1

$$\bar{x}' = \begin{pmatrix} 3 & -2 \\ 2 & -2 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 3x_1 - 2x_2 \\ 2x_1 - 2x_2 \end{pmatrix}$$

I) Try $\bar{x} = e^{rt} \bar{\epsilon}$. r must be an eigenvalue to A and $\bar{\epsilon}$ the corresponding eigenvector. See book

II) Change to new coordinates (functions)
 $\begin{pmatrix} y_1(t) \\ y_2(t) \end{pmatrix}$ in such a way that we get two 1D ODE problems. That is

$$\bar{x} = T \bar{y}$$



$$x'_1 = 3x_1 - 2x_2$$

$$x'_2 = 2x_1 - 2x_2$$

$$x'_1 = 0 \text{ when } x_2 = \frac{3x_1}{2}$$

$$x'_2 = 0 \text{ when } x_2 = x_1$$

(2) in the γ_1 direction

(1) in the γ_2 direction

$$\begin{cases} \gamma'_1 = +2\gamma_1 \\ \gamma'_2 = -\gamma_2 \end{cases}$$

Inward motion along γ_2 -axis
Outward motion along γ_1 -axis

$$\bar{x}' = T\bar{y}' = A\bar{x} = AT\bar{y}$$

$$T\bar{y}' = AT\bar{y}$$

$$\bar{y}' = \underbrace{T^{-1}AT}_{D}\bar{y}$$

We want D to be a diagonal matrix. Then the columns of T must be the eigenvectors of A .

$$\begin{pmatrix} 3 & -2 \\ 2 & -2 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \lambda \begin{pmatrix} c_1 \\ c_2 \end{pmatrix}$$

Non-trivial solution $\Leftrightarrow \begin{vmatrix} 3-2 & -2 \\ 2 & -2-\lambda \end{vmatrix} = 0$

$\lambda_1 = 2$ has eigenvector $\begin{pmatrix} 2 \\ 1 \end{pmatrix}$. $\lambda_2 = -1$ has eigenvector $\begin{pmatrix} 1 \\ 2 \end{pmatrix}$.

$$D = \begin{pmatrix} 2 & 0 \\ 0 & -1 \end{pmatrix}$$

Easy to solve $y_1' = 2y_1$,
 $y_2' = -y_2$

The equations don't "talk" to each other.

$$y_1 = C_1 e^{2t}, y_2 = C_2 e^{-t}$$

Back to \bar{x}

$$\bar{x} = \begin{pmatrix} 2 & 1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} C_1 e^{2t} \\ C_2 e^{-t} \end{pmatrix}$$

$$\bar{x} = C_1 e^{2t} \begin{pmatrix} 2 \\ 1 \end{pmatrix} + C_2 e^{-t} \begin{pmatrix} 1 \\ 2 \end{pmatrix}$$

Expansion in $\begin{pmatrix} 2 \\ 1 \end{pmatrix}$ direction

Contraction in $\begin{pmatrix} 1 \\ 2 \end{pmatrix}$ direction

The CP \bar{o} is here
a saddle point.

What happens if the eigenvalues
are complex? See 7.6

What happens when we have
less than n eigenvectors? See 7.8

So far
we have

$$\begin{array}{c} \text{---} \\ | \quad o \quad o \end{array} \rightarrow \lambda \quad \bar{o} \text{ unstable}$$

$$0 \quad \lambda_1 \quad \lambda_2$$

3 cases:

$$\begin{array}{c} \text{---} \\ | \quad \lambda_1 \quad \lambda_2 \\ o \end{array} \rightarrow \lambda \quad \bar{o} \text{ saddle point}$$

$$\begin{array}{c} \text{---} \\ | \quad + \quad + \\ \lambda_1 \quad \lambda_2 \quad o \end{array} \rightarrow \lambda \quad \bar{o} \text{ asymptotically stable.}$$

Ex 7.5, 11) $\bar{x}' = \begin{pmatrix} 1 & 1 & 2 \\ 1 & 2 & 1 \\ 2 & 1 & 1 \end{pmatrix} \bar{x}$

Diagonalisation
leads to the eq.

$$\begin{vmatrix} 1-\lambda & 1 & 2 \\ 1 & 2-\lambda & 1 \\ 2 & 1 & 1-\lambda \end{vmatrix} = 0$$

cubic equation

Note 4 is an
eigenvalue with eigenvector $\begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$. Then
easy to find the other eigenvalues.

$$\bar{x} = C_1 e^{4t} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} + C_2 \begin{pmatrix} 1 \\ -2 \\ 1 \end{pmatrix} e^t + C_3 \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} e^{-t}$$

Ex 7.6.1)

$$A = \begin{pmatrix} 3 & -2 \\ 4 & -1 \end{pmatrix}$$

$$\begin{vmatrix} 3-\lambda & -2 \\ 4 & -1-\lambda \end{vmatrix} = 0$$

$$\lambda = 1 + i2$$

We denote the eigenvector to $\lambda = 1 + 2i$
 $\bar{a} + i\bar{b}$. One solution to ODE

$$\bar{x}' = A\bar{x} \text{ is}$$

$$e^{(1+2i)t} (\bar{a} + i\bar{b}) =$$

$$e^t [(\bar{a}\cos 2t - \bar{b}\sin 2t) + i(\bar{a}\sin 2t + \bar{b}\cos 2t)]$$

Real-valued solutions to ODE
are

$$\bar{U}(t) = e^t (\bar{a}\cos 2t - \bar{b}\sin 2t)$$

and

$$\bar{V}(t) = e^t (\bar{a}\sin 2t + \bar{b}\cos 2t)$$

$i \rightarrow -i$ gives no new solutions.

What is \bar{a} and \bar{b} ?

$$\begin{pmatrix} 3 & -2 \\ 4 & -1 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = (1+2i) \begin{pmatrix} c_1 \\ c_2 \end{pmatrix}$$

$$\begin{pmatrix} 1 \\ 1-i \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} + i \begin{pmatrix} 0 \\ -1 \end{pmatrix}, \quad \bar{\alpha} = \begin{pmatrix} 1 \\ i \end{pmatrix}$$

$$\bar{b} = \begin{pmatrix} 0 \\ -1 \end{pmatrix}. \quad \text{Finally}$$

$$\bar{v}(t) = e^t \begin{pmatrix} \cos 2t \\ \cos 2t + \sin 2t \end{pmatrix}$$

$$\bar{r}(t) = e^t \begin{pmatrix} \sin 2t \\ \sin 2t - \cos 2t \end{pmatrix}$$

\bar{o} is a spiral point and it is asymptotically unstable.

For example if $\bar{x}(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix} = C_1 \bar{v}(0) + C_2 \bar{r}(0)$ we get

$$C_1 = C_2 = 1$$

$$\bar{x}(t) = e^t \begin{pmatrix} \cos 2t + \sin 2t \\ 2 \sin 2t \end{pmatrix}$$

see next page

$$\text{Ex 7.6.28} \quad m v''(t) + k v(t) = 0$$

$$x_1 = v$$

$$x_2 = v'$$

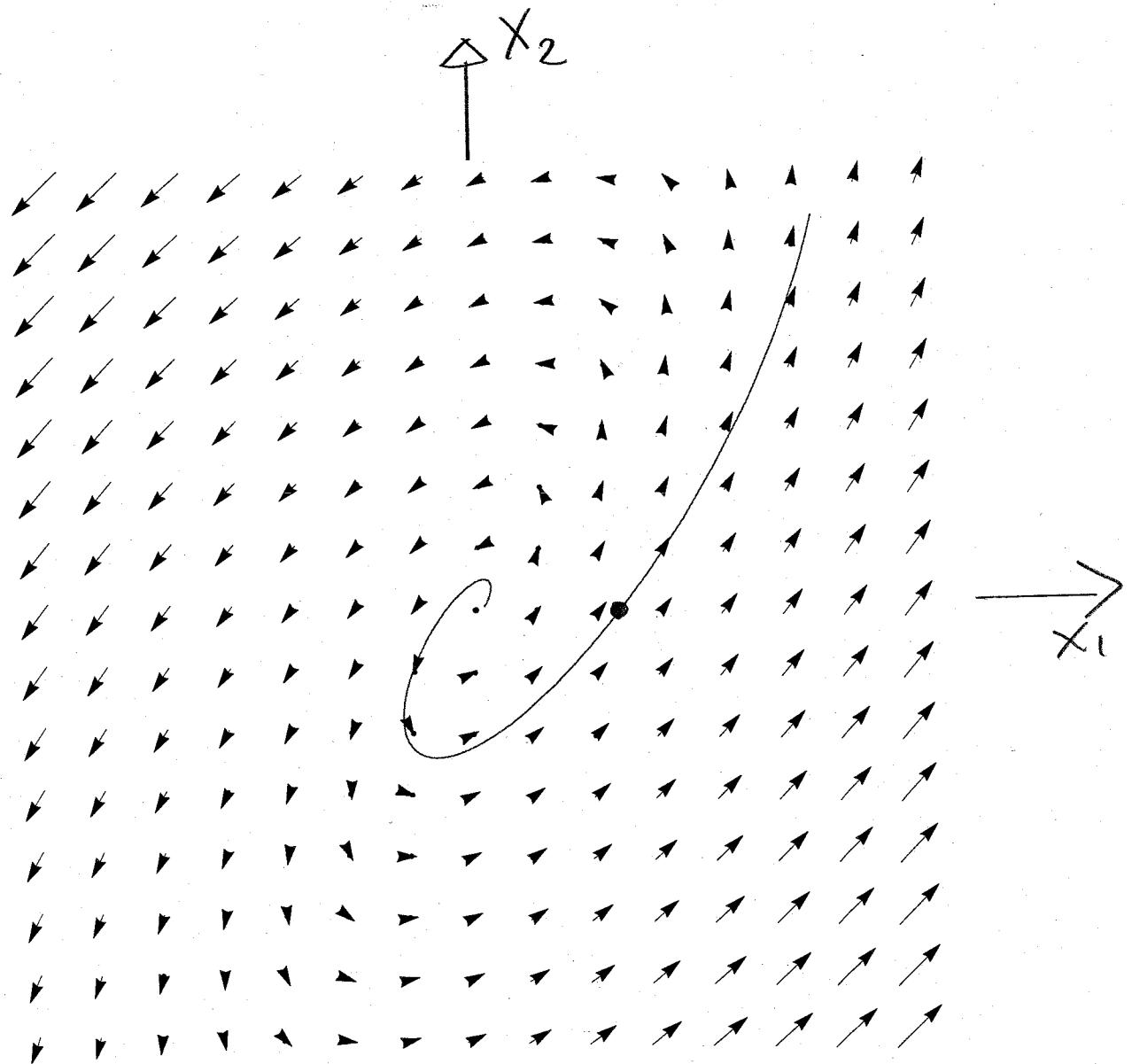
Spring
Without
friction

$$x_1' = x_2$$

$$x_2' = v'' = \frac{-k}{m} v = -\frac{k}{m} x_1$$

$$A = \begin{bmatrix} 0 & 1 \\ -\frac{k}{m} & 0 \end{bmatrix}$$

Out[36]=



$$\bar{x}(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$-3 \leq t \leq \frac{1}{2}$$

$$\begin{cases} x_1' = 3x_1 - 2x_2 \\ x_2' = 4x_1 - x_2 \end{cases}$$

PUT $\frac{k}{m} = \omega^2$. Solving

$$\begin{vmatrix} -\lambda & 1 \\ -\omega^2 & -\lambda \end{vmatrix} = 0 \Leftrightarrow$$

$$\lambda^2 + \omega^2 = 0 \Leftrightarrow \lambda = \pm i\omega$$

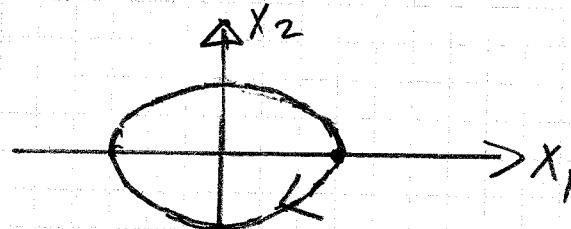
$$\bar{a} + i\bar{b} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} + i \begin{pmatrix} 0 \\ \omega \end{pmatrix} \text{ check!}$$

$$\bar{U}(t) = \cos\omega t \begin{pmatrix} 1 \\ 0 \end{pmatrix} - \sin\omega t \begin{pmatrix} 0 \\ \omega \end{pmatrix} = \begin{pmatrix} \cos\omega t \\ -\omega \sin\omega t \end{pmatrix}$$

$$\bar{V}(t) = \sin\omega t \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \cos\omega t \begin{pmatrix} 0 \\ \omega \end{pmatrix} = \begin{pmatrix} \sin\omega t \\ \omega \cos\omega t \end{pmatrix}$$

$|\bar{U} \times \bar{V}| = \omega$, so \bar{U} and \bar{V} are linearly independent for all times.

If $\bar{x}(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ then $\bar{x}(t) = \bar{U}(t)$



\bar{o} is stable.

$n=2$

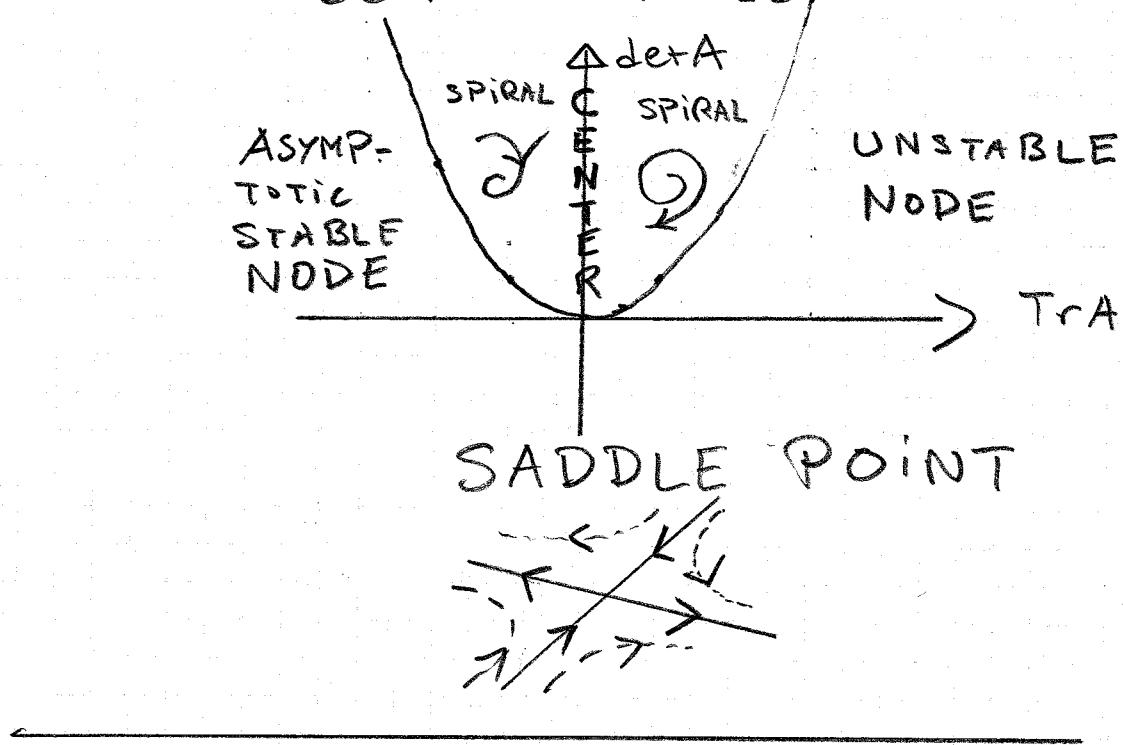
overview

$$\begin{vmatrix} a_{11} - \lambda & a_{12} \\ a_{21} & a_{22} - \lambda \end{vmatrix} = 0$$

$$\lambda_{1,2} = \frac{\text{Tr } A}{2} \pm \sqrt{\left(\text{Tr } A\right)^2 - 4 \det A}$$

Here $\text{Tr} A = \alpha_{11} + \alpha_{22}$ and

$$\det A = \alpha_{11} \cdot \alpha_{22} - \alpha_{21} \cdot \alpha_{12}$$



On next page you see some illustrations of what can happen in $n=3$ case.

From V. I. Arnold
 "ODE"

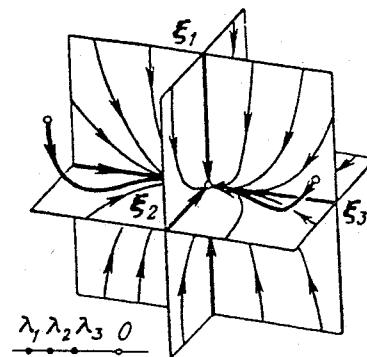


Fig. 137 Phase space of a linear equation in the case $\lambda_1 < \lambda_2 < \lambda_3 < 0$. The phase flow is a contraction in all three directions.

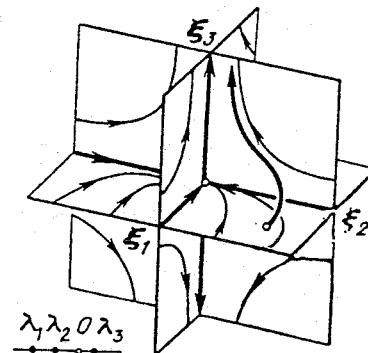


Fig. 138 The case $\lambda_1 < \lambda_2 < 0 < \lambda_3$: Contraction in two directions and expansion in the third.

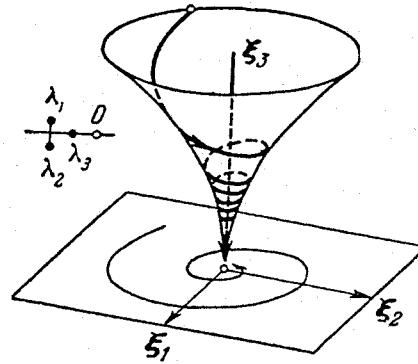


Fig. 139 The case $\operatorname{Re} \lambda_{1,2} < \lambda_3 < 0$: Contraction in the direction of ξ_3 and rotation with faster contraction in the plane of ξ_1 and ξ_2 .

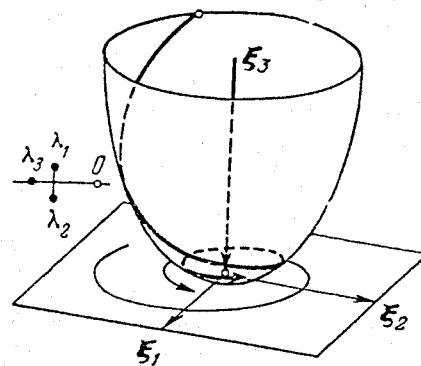


Fig. 140 The case $\lambda_3 < \operatorname{Re} \lambda_{1,2} < 0$: Contraction in the direction of ξ_3 and rotation with slower contraction in the plane of ξ_1 and ξ_2 .

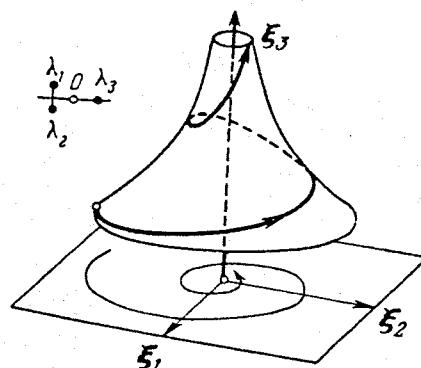


Fig. 141 The case $\operatorname{Re} \lambda_{1,2} < 0 < \lambda_3$: Expansion in the direction of ξ_3 and rotation with contraction in the plane of ξ_1 and ξ_2 .

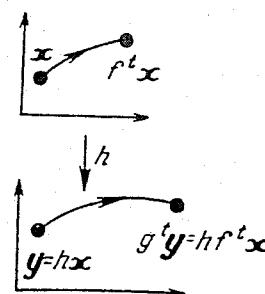


Fig. 142 Equivalent flows.